Zeta Function Geometric approximation **Ergodic Tools** Locating zeros Bibliography Illusions: curves of xeros of Felberg xeta functions

Polina Vytnova joint work with Mark Pollicott

University of Warwick

On one property of one analytic function

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Selberg Zeta Function

Let X be a compact surface of constant negative sectional curvature $\kappa = -1$. Define

$$Z_X(s) = \prod_{n=0}^{\infty} \prod_{\substack{\gamma = ext{primitive} \ ext{closed geodesic}}} \left(1 - e^{-(s+n)\ell(\gamma)}
ight),$$

Theorem (Selberg, 1956)

Let X be a compact Riemann surface. Then the infinite product converges to an analytic non-zero function on $\Re(s) > 1$ and extends as an analytic function to \mathbb{C} . The function Z_X has a simple zero at s = 1 and for any zero s in the critical strip $0 < \Re(s) < 1$ we have that either $s \in [0, 1]$ is real, or $\Re(s) = \frac{1}{2}$.

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Years Past...





ENIAC and its first programmers, c.1950

Dell Mini, 2017

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Numerical Experiments Revealed



Figure: 29504 Zeros of *an approximation* to the Selberg zeta function associated to a pair of pants. D. Borthwick, 2014

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This Plot Raised Many Questions

- What exactly the approximation is? (An infinite product can't be evaluated numerically, unless it can be reduced to a finite one.)
- If we consider another approximation to the same function, will the plot be different?
- (3) Are these zeros any close to the zeros of ζ ?
- Why do we see the curves?
- 5 If we consider another surface, how the plot will change?

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Another Example



Figure: 107164 Zeros of the Selberg zeta function associated to a one-holed torus. P.V., 2018.

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Why Do We See the Curves?

It is a feature (or a bug) of the outlook we have, like the photo below.



Figure: P.V. holding the Hunter's moon on the 24th of October.

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Disappearance of the Curves

Take an affine transform for a closer look



Figure: A zoom-in of the plot of the zero set of the Selberg's zeta for a pair of pants.

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A One-Holed Torus



- Topologically one-holed torus *T* is a punctured sphere with a handle;
- It is a surface of constant negative curvature −1 and cannot be embedded into ℝ³ by Efimov's theorem;
- As a metric space, it is uniquelly defined by the lengths of two geodesics and the angle inbetween $T = T(\ell_1, \ell_2, \varphi)$;
- It possess countably many closed geodesics $\{\gamma_n\}$ of lengths $0 < \ell(\gamma_1) < \ell(\gamma_2) < \ldots < \ell(\gamma_n) \ldots \to \infty$
- Symmetric torus means $\ell_1 = \ell_2$, $\varphi = \frac{\pi}{2}$.

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A Pair of Pants

Surfaces

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- Topologically pair of pants X is a 3-punctured sphere;
 - It is a surface of constant negative curvature -1 and cannot be embedded into \mathbb{R}^3 by Efimov's theorem;
- As a metric space, it is uniquelly defined by the lengths of the three boundary geodesics: $X = X(\ell_1, \ell_2, \ell_3)$;
- It possess countably many closed geodesics $\{\gamma_n\}$ of lengths $0 < \ell(\gamma_1) < \ell(\gamma_2) < \ldots < \ell(\gamma_n) \ldots \to \infty$
- Symmetric pair of pants means $\ell_1 = \ell_2 = \ell_3 =: b.$

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The Hyperbolic Action



 Cutting the pair of pants along the red geodesics, we obtain a pair of hexagons;

The hexangons can be immersed into \mathbb{H}^2 as right-angled hexagons;

- The Fuchsian group Γ = ⟨R₁, R₂, R₃⟩, generated by reflections with respect to the "cuts", gives a pair of pants as a double cover of the factor space X(b) = ℍ²/Γ;
- To any geodesic X corresponds a geodesic in ℍ; for any closed geodesic γ there exists R_γ ∈ Γ preserving γ.
- The action $\Gamma \curvearrowright \mathbb{H}^2$ is hyperbolic.

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Properties of Selberg Zeta Functions

- In 1992, Guillopé established that in the case of geometrically finite hyperbolic surfaces of infinite area, the infinite product Z_X converges for ℜ(s) sufficiently large and has a meromorphic extension to C.
- 2 Zeros of the Selberg zeta function correspond to the poles of the Ruelle zeta function given by

$$\zeta(s)$$
: $=rac{Z_X(s+1)}{Z_X(s)} = \prod_{\substack{\gamma = ext{primitive} \ ext{closed geodesic}}} \left(1 - e^{-s\ell(\gamma)}
ight)^{-1}$

- 3 There exists the largest real zero δ, which is equal to the Hausdorff dimension of the limit set of Γ (a subset of the unit circle).
- ④ There is no other zeros with $\Re(s) = \delta$



Properties of Selberg Zeta Functions (continued)

- δ is the growth rate of the number of primitive closed geodesics δ = lim_{t→∞} ¹/_t log #{γ: ℓ(γ) ≤ t}. Moreover, #{γ: ℓ(γ) ≤ t} ~ ^{e^{δt}}/_{δt}.
- For a symmetric pair of pants $\delta = \delta(b) \sim \frac{1}{b}$ (McMullen)
- 7 There exists $\varepsilon > 0$ such that there is only finite number of zeros satisfying $\Re(s) > \delta \varepsilon$ (Jakobson–Naud)
- Complex zeros are related to the eigenvalues of the Laplacian operator acting on L₂ functions and are a subject of intensive research (Nonnenmacher, Patterson, Perry, Zworski ...). These are defined as the poles of the resolvent and are referred to as *resonances* of X.

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Closed Geodesics



To every closed geodesic γ on X(b) corresponds

• a cutting sequence of period 2n

 $\cdots j_{2n-1}j_{2n}j_{2n+1}\cdots$

where $j_k \in \{1, 2, 3\}$, $j_k \neq j_{k+1}$ for $1 \leq k \leq 2n$ and $j_{2n} \neq j_1$.

a periodic orbit of the subshift σ of finite type on the space of 3 symbols $\Sigma = \{1, 2, 3\}^{\mathbb{Z}}$ with transition matrix

$$\begin{pmatrix} 0 & 1 & 1 \\ 1 & 0 & 1 \\ 1 & 1 & 0 \end{pmatrix}$$

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Transition Matrices

Let's fix *n* and define $r_n \colon \Sigma \to \mathbb{R}$, $r_n(\xi) = \ell(\gamma_{[\xi_{[n/2]}, \xi_{[n/2]+1}}])$, where γ is chosen such that

$$\ell(\gamma) = \min_{\gamma'} \{\ell(\gamma') \mid \gamma' \text{ intersects } \xi_1, \dots, \xi_n\}$$

Let ξ^1, \ldots, ξ^N be all subsequences of the sequences in Σ of the length *n*. We define an $N \times N$ transition matrix

$$A_{i,j}^n = \begin{cases} 1, & \text{if } \xi_{k+1}^i = \xi_k^j; \text{ for } k = 1, \dots, n-1 \\ 0, & \text{otherwise.} \end{cases}$$

and a complex matrix function

$$A: \mathbb{C} \to Mat(N, N) \qquad A_{i,j}(s) = \exp(-sr_n(\xi)) \cdot A_{i,j}^n,$$

where $\xi = \xi_1^i \dots \xi_n^i \xi_n^j$.

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Key Lemma

Lemma

$$\left[\left(1 - e^{-s\ell(\gamma)} \right)^2 = \lim_{n o \infty} \det \left(I_N - A^2(s) \right);
ight.$$

 $\gamma = primitive$ closed geodesic

where I_N is the $N \times N$ identity matrix.

Choosing
$$n=2$$
 above we get $r_2\equiv b$
 $\det(\mathit{Id}-e^{-2sb}A^2)=(1-4e^{-2bs})(1-e^{-2bs})^2$

For a first approximation...

- The zero set belongs to a pair of straight lines
- The distance between consequetive zeros is $\frac{\pi}{h}$.

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Using n = 3 in the approximation of geodesics length

$$r_3(\xi) = b + c(\xi)e^{-b} + O(e^{-2b}),$$

we obtain a 6×6 matrix which determinant has the zero set on the curves

$$\begin{split} \mathcal{C}_{1} &= \left\{ \frac{1}{2b} \ln |2 - 2\cos(t)| + ie^{b}t \mid t \in \mathbb{R} \right\}; \\ \mathcal{C}_{2} &= \left\{ \frac{1}{2b} \ln |2 + \cos(2t)| + ie^{b}t \mid t \in \mathbb{R} \right\}; \\ \mathcal{C}_{3} &= \left\{ \frac{1}{2b} \ln \left| 1 - \frac{1}{2}e^{2it} - \frac{1}{2}e^{it}\sqrt{4 - 3e^{2it}} \right| + ie^{b}t \mid t \in \mathbb{R} \right\}; \\ \mathcal{C}_{4} &= \left\{ \frac{1}{2b} \ln \left| 1 - \frac{1}{2}e^{2it} + \frac{1}{2}e^{it}\sqrt{4 - 3e^{2it}} \right| + ie^{b}t \mid t \in \mathbb{R} \right\}. \end{split}$$

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Curves of Zeros — II



Figure: The zero sets of $\zeta_X(\frac{\sigma}{b} + ite^b)$ and renormalized curves C_k , for b = 6; and a zoomed neighbourhood of $(\frac{\ln 2}{2}, \frac{\pi}{4})$.

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Comments on Geometric Approximation

- 1 Increasing *n* we do not see a change in the zero set for $\Im(z) < e^{3b}$;
- 2 There is no good estimates on error term (or rate of convergence).

We need to estimate the approximation error.

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Transfer Operators Technique

Given a hyperbolic action, we introduce:

- A proper Banach space of analytic functions;
- 2 A nuclear transfer operator acting on the Banach space;
- 3 The determinant of the transfer operator, which is an analytic function;
- Ruelle–Pollicott dynamical zeta function;
- The Ruelle zeta function turns to be an analytic function, which is closely related to the determinant (of the transfer operator);
- The zeta function can be computed very efficiently using periodic orbits data (of the hyperbolic system) and its zeros provide quontitative information about the system.

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The space \mathcal{B} of analytic functions on the union of disjoint disks $\bigsqcup_{k=1}^{3} U_k$, chosen so that $R_i(U_j \cup U_k) \subset U_i$ for any three distinct $i, j, k \in \{1, 2, 3\}$.



Figure: The domain of analytic functions forming the Banach space (in pale red).

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Transfer Operator

We define a transfer operator \mathcal{L}_{s} on the space $\mathcal B$ by

$$(\mathcal{L}_{s}f)|_{U_{1}}(z_{1}) = |R'_{1}(z_{2})|^{s}f(z_{2}) + |R'_{1}(z_{3})|^{s}f(z_{3}),$$

where z_2, z_3 are preimages of $z_1 \in U_1$ with respect to reflection with respect to the geodesic β_1 .

Lemma (Grothendieck–Ruelle) The operator \mathcal{L}_{s} is nuclear.

We may write the determinant of the transfer operator as

$$\zeta(z,s) \stackrel{\text{def}}{=} \exp\left(-\sum_{n=1}^{\infty} \frac{z^n}{n} \operatorname{Tr} \mathcal{L}_s^n\right) = \det(I - z\mathcal{L}_s).$$

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Zeta Function Magic

Lemma (Grothendieck-Ruelle)

The trace of the transfer operator may be explicitly computed in terms of the closed geodesics.

$$\operatorname{Tr} \mathcal{L}_{s}^{n} = \sum_{|\gamma|=n} \frac{\exp(-s\ell(\gamma))}{1 - \exp(-\ell(\gamma))}$$

Theorem (Ruelle)

There exists a constant δ such that the determinant is an analytic function in both variables in a strip $0 < s < \delta$, and

$$\zeta(1,s) = \zeta(s) = \exp\Bigl(\sum_{n=1}^{\infty} \frac{1}{n} \sum_{|\gamma|=n} \frac{\exp(-s\ell(\gamma))}{1 - \exp(-\ell(\gamma))}\Bigr)$$

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Computing the Zeta Function

Using Ruelle's Theorem,

$$\zeta(s) = \sum_{n=0}^{\infty} z^n a_n(s) \Big|_{z=1} = \lim_{N \to \infty} \sum_{n=0}^{N} a_n(s),$$

where a_n are explicitly defined in terms of closed geodesics of the word length not more than $|\gamma| \leq 2n$, and are analytic in s:

$$\mathsf{a}_n(s) = -rac{1}{n}\sum_{j=0}^{n-2} \mathsf{a}_j(s) \cdot \mathrm{Tr}\mathcal{L}_s^{n-j}$$

Lemma (after Grothendieck–Ruelle) The terms $a_n(s)$ are decreasing superexponentially: $|a_n(s)| < \lambda(s)^{n^2}$, where $\lambda(s) < 1$ depend only on \mathcal{L}_s , but the estimate is not uniform in s.

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Algorithm

Choosing truncation $\zeta_N(s) = \sum_{n=0}^N a_n(s)$, we can

- Ind the largest real zero = the width of the critical strip,
- 2 consider a dense lattice in the strip,
- ③ compute the residue over each square,
- ④ find a zero using Newton method starting from a point of the lattice.

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Numerical Output: Symmetric Pants



Figure: Zeros of the zeta function associated to a symmetric pair of pants and a more careful look for b = 12, N = 14.

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Another Viewpoint: Exponential Sums

The function $\zeta_N(s)$ is a finite exponential sum

$$\zeta_N(s) = \sum_{j=k}^n \alpha_k \exp(\mu_k s),$$

where the multipliers μ_k are the lengths of closed geodesics with word length up to 2N.

- Zeros form a point-periodic set and belong to a finite strip, parallel to the imaginary axis
- 2 Their imaginary parts satisfy relation

$$\Im(s_k) = \frac{\pi}{\max \mu_k - \min \mu_k} + \varphi(k),$$

for an almost periodic function φ .

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Main Approximation Result

$$\mathcal{R}(\mathcal{T}) = \{ s \in \mathbb{C} \mid 0 \leq |\Re(s)| \leq \delta \text{ and } |\Im(s)| \leq \mathcal{T} \}.$$

Theorem (M. Pollicott-P. V.)

Let X be a symmetric pair of pants with boundary geodesics of the length $\ell(\gamma_0) = 2b$. We may approximate ζ on the domain $\mathcal{R}(T)$ by a complex trigonometric polynomial ζ_n so that $\sup_{\mathcal{R}(T)} |\zeta - \zeta_n| \le \eta(b, n, T)$, where $T(b) = e^{k_0 b}$ for some constant $1 < k_0 < 2$ independent of b and n, such that

- 1 for any $n \ge 14$ we have $\eta(b, n, T(b)) \le O\left(\frac{1}{\sqrt{b}}\right)$ as $b \to \infty$
- 2 for any $b \ge 20$ we have $\eta(b, n, T(b)) \le O(e^{-bk_1n^2})$ as $n \to \infty$.

for some $k_1 > 0$ which is independent on b and n.

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Subsequent Approximations



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Final Approximation

Lemma (after M. Pollicott-P.V.)

There exists an explicit 6-by-6 matrix B(s) such that the real analytic function $\zeta_{12}(\frac{\sigma}{b} + ite^{b})$ converges uniformly to $\det(I - e^{-2\sigma - 2itbe^{b}}B(e^{it}))$, and more precisely,

$$\left|Z_{12}\left(\frac{\sigma}{b}+ite^{b}\right)-\det\left(I-e^{-2\sigma-2itbe^{b}}B(e^{it})\right)\right|=O\left(e^{-b}\right)$$

as $b \to +\infty$.

- The matrix B can be constructed using a transition matrix of a subshift of finite type on the space {1, 2, 3}^ℕ.
- \bullet The curves $\mathcal{C}_1, \mathcal{C}_2, \mathcal{C}_3, \mathcal{C}_4$ computed using the formula

$$|e^{2\sigma}| = eig(B(e^{it}))$$

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Thank you!